

# Local simulation of singlet statistics for restricted set of measurement

Joint Estonian-Latvian Theory Days  
(2nd-5th Oct. 2014)

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# Quantum State of physical systems ...

Pure State:  $|\psi\rangle \in \mathbb{C}^n$

- ▶ Qubit:  $|\psi\rangle = a|0\rangle + b|1\rangle \in \mathbb{C}^2$ ;  
 $a, b \in \mathbb{C}; |a|^2 + |b|^2 = 1$ .
- ▶ Two-qubits:  $|\psi\rangle = a|00\rangle + b|01\rangle + c|10\rangle + d|11\rangle \in \mathbb{C}^4$ ;  
 $a, b, c, d \in \mathbb{C}, |a|^2 + |b|^2 + |c|^2 + |d|^2 = 1$ .

Mixed State  $\in \mathbb{C}^n$

- ▶ A mixture of pure states  $\{|\psi_i\rangle, p_i\}$  where  $|\psi_i\rangle \in \mathbb{C}^n$
- ▶ **Density matrix:** A more compact representation is  $\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|$ , where  $\text{Tr}(\rho)=1$  and  $\rho \geq 0$ .
- ▶ A mixed state  $\rho$  can be prepared in many different ways but physically they are all same.

## Quantum Measurement / Quantum Observables)...

- ▶ Outcome of measurement on a quantum system in general is probabilistic.
- ▶ Quantum measurement on any state  $|\psi\rangle$  can be expressed as an Hermitian operator  $\hat{O}$ .
- ▶ Possible results of measurement are eigenvalues  $i$  of  $\hat{O}$ .
- ▶ Post measurement state after getting an outcome  $i$  is  $|i\rangle$ : the corresponding eigenvectors of  $\hat{O}$ .
- ▶ **Born-Rule:** Probability of getting  $i$ -th outcome is given by  $P(i) = |\langle\psi|i\rangle|^2$ .
- ▶ Example: Measuring  $\hat{O} = |0\rangle\langle 0| - |1\rangle\langle 1|$  on a single qubit state  $a|0\rangle + b|1\rangle$  results in two possible outcome,  $+1$  with probability  $|a|^2$  and  $-1$  with probability  $|b|^2$ ; corresponding post measurement state is  $|0\rangle$  and  $|1\rangle$  respectively.

# Quantum Entanglement

In quantum mechanics state of a composite system constituting of two or more subsystems can be entangled.

## Consider a two qubit bi-partite system

- ▶ One qubit is with Alice and other qubit is with Bob.
- ▶ Then state of composite system  $|\psi_{AB}\rangle \in \mathbb{C}^2 \otimes \mathbb{C}^2$ .
- ▶ If  $|\psi_{AB}\rangle \neq |\psi_A\rangle \otimes |\psi_B\rangle$  for any  $|\psi_A\rangle, |\psi_B\rangle \in \mathbb{C}^2$  then such states are called entangled states.
- ▶ **Singlet state:**  $|\psi^-\rangle = \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle)$  is an entangled state.

## A two-qubit mixed entangled state

**Werner-State:** Is a mixture of  $|\psi^-\rangle$  and White noise  $\frac{\mathbb{I}}{4}$   
 $\rho_W = p|\psi^-\rangle\langle\psi^-| + (1-p)\frac{\mathbb{I}}{4}$  where  $0 \leq p \leq 1$

Werner states are entangled for  $p > \frac{1}{3}$

# Entanglement and Nonlocality

Suppose Alice and Bob, located far apart, share Singlet state  $|\psi^-\rangle$ . Now, if they measure their respective subsystems, then their local measurement outcomes can be **non-classically correlated**. This can also be true for many other entangled states.

## Non-classical Correlations

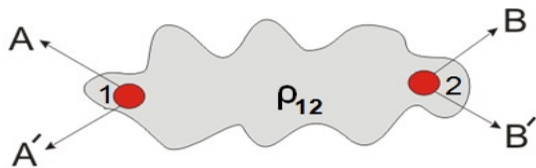
- ▶ Let Alice's (Bob's) measurement be  $A$  ( $B$ ) with possible outcomes  $a(b) \in \{-1, +1\}$ .
- ▶ Quantum Mechanics:  $P_Q(a, b|A, B) = \text{Tr}(|\psi^-\rangle\langle\psi^-|A \otimes B)$ .
- ▶ **Local Hidden Variables (LHV)**: If Alice and Bob by pre-sharing some local variable  $\lambda_i$  and without later communication can simulate Quantum probabilities, then we get a LHV model

$$P_Q(a, b|A, B) = \int_{\Lambda} d\lambda D(\lambda) P(a|A, \lambda) P(b|B, \lambda)$$

## A Nonlocality Witness

Suppose Alice (Bob) randomly choose to measure  $\{A, A'\}$  ( $\{B, B'\}$ ) and possible outcome of their measurement  $\in \{-1, +1\}$ . Then if there is some LHV model for their measurement statistics then the following constraint must be satisfied:

Bell-CHSH inequality



$$-2 \leq \langle AB \rangle + \langle A'B \rangle + \langle AB' \rangle - \langle A'B' \rangle \leq 2$$

# Motivation...

## Entanglement $\neq$ Nonlocality

- ▶ Pure singlet state show maximum Bell-violation for appropriately chosen ideal projective measurements.
- ▶ Can all entangled states generate some non-classical (nonlocal) correlation?
- ▶ **Werner's result**—Projective measurements on mixed entangled states  $\rho_{AB} = p|\psi^-\rangle\langle\psi^-| + (1-p)\frac{\mathbb{I}}{4}$  where  $\frac{1}{3} < p \leq \frac{1}{2}$  cannot generate nonlocal correlation.
- ▶ Werner showed that, above mixed (impure) entangled states can be locally simulated by two spatially separated parties by pre-sharing classical correlations (Local Hidden Variables).
- ▶ We have studied this problem from the opposite direction i.e. rather than weakening the state we restrict the class of observable to provide local model for the pure singlet state.

## General quantum observables

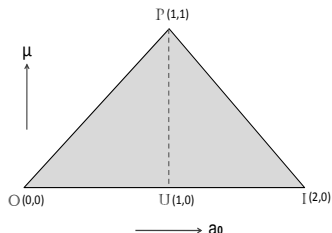
- ▶ Generalized quantum observables are described by positive operator valued measure (POVM).
- ▶ POVM is a collection of selfadjoint operators  $\{E_i\}$  acting on state space such that:
  - (i)  $0 \leq E_i \leq I$  for all  $i$ ,
  - (ii)  $\sum_i E_i = I$ , where  $i \in \{1, 2, \dots, n\}$ .
- ▶ Measurement  $\{E_i\}$  on a quantum state  $\rho$  results in any one of the  $n$  possible outcomes; probability of occurrence of  $i$ -th outcome (termed as clicking of  $i$ -th effect) is  $Tr[\rho E_i]$ .



## General two-outcome measurement on a qubit

Let unit vector  $\hat{a} = (a_1, a_2, a_3) \in \mathbb{R}^3$  and  $\hat{a} \cdot \vec{\sigma} = a_1\sigma_1 + a_2\sigma_2 + a_3\sigma_3$  where  $\sigma_i$  are Pauli matrices:  $\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$ ,  $\sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$ ,  $\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

POVM  $\{E, I-E\}$  acting on  $\mathbb{C}^2$



$$E = \frac{1}{2}[a_0 I + \mu \hat{a} \cdot \vec{\sigma}] \quad (1)$$

$$0 \leq a_0 \leq 2 \quad (2)$$

$$0 \leq \mu \leq \min\{a_0, 2 - a_0\} \quad (3)$$

- ▶ Point  $P(1,1)$  correspond to an ideal projective measurement.
- ▶ Measurements corresponding to points on the dashed line  $UP$  can be physically interpreted as unsharp-spin property of a spin- $\frac{1}{2}$  system (P. Busch, 1986).

# Singlet statistics for general two-outcome measurements

- ▶ Suppose, two spatially separated parties Alice and Bob share one qubit each from a singlet state

$$|\psi^-\rangle = \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle)$$

- ▶ Let Alice's (Bob's) observable be a most general two-outcome POVM  $E_A[a_0, \mu_A, \hat{a}]$  ( $E_B[b_0, \mu_B, \hat{b}]$ )

## Joint outcome probabilities

$$P^{AB}(\text{yes}, \text{yes}) = \frac{1}{4}[a_0 b_0 - \mu_A \mu_B \hat{a} \cdot \hat{b}]$$

$$P^{AB}(\text{yes}, \text{no}) = \frac{1}{4}[a_0(2 - b_0) + \mu_A \mu_B \hat{a} \cdot \hat{b}]$$

$$P^{AB}(\text{no}, \text{yes}) = \frac{1}{4}[(2 - a_0)b_0 + \mu_A \mu_B \hat{a} \cdot \hat{b}]$$

$$P^{AB}(\text{no}, \text{no}) = \frac{1}{4}[(2 - a_0)(2 - b_0) - \mu_A \mu_B \hat{a} \cdot \hat{b}]$$

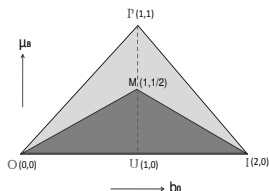
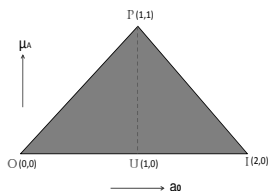
## LHV models for singlet statistics

- ▶ There can be no local hidden variable model for pure singlet state statistics generated by ideal projective measurements (for suitable choice of measurement directions Bell-CHSH inequality is violated)
- ▶ Our motivation is to explore possibilities of local hidden variable model for the pure singlet state by restricting (deviating from ideal projective measurements) the parameters of a two-outcome POVM measurement.
- ▶ We give two forms of LHV models for singlet state under certain restrictions on parameters of two outcome POVMs.
- ▶ In both the models vectors  $\hat{\lambda} = (\sin \theta \cos \phi, \sin \theta \sin \phi, \cos \theta)$  uniformly distributed over unit sphere are local variables pre-shared between Alice and Bob.

# A fully biased model $M_{fb}$

## Restriction on observable

- ▶ Alice: No restriction
- ▶ Bob:  $\mu_B \leq \frac{1}{2} \min\{b_0, 2 - b_0\}$



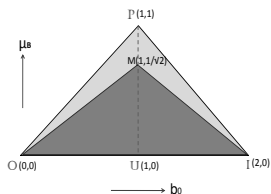
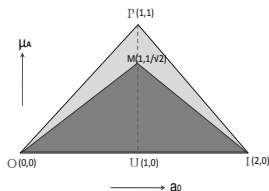
## Model

- ▶ Alice:  $P_{\hat{\lambda}}^A(\text{yes}) = \frac{a_0}{2} + \frac{1}{2} \mu_A \cos \alpha$
- ▶ Bob:  $P_{\hat{\lambda}}^B(\text{yes}) = \frac{b_0}{2} - \mu_B \operatorname{sgn}(\cos \beta)$
- ▶  $\alpha$  ( $\beta$ ) is angle with pre-shared variable  $\hat{\lambda}$
- ▶  $\operatorname{sgn}(x) = +1$  ( $-1$ ) for  $x \geq 0$  ( $x < 0$ )

# A fully symmetric model $\mathbb{M}_{fs}$

## Restriction on observable

- ▶ Alice:  $\mu_A \leq \frac{1}{\sqrt{2}} \min\{a_0, 2 - a_0\}$
- ▶ Bob:  $\mu_B \leq \frac{1}{\sqrt{2}} \min\{b_0, 2 - b_0\}$



## Model

- ▶ Alice:  $P_{\hat{\lambda}}^A(\text{yes}) = \frac{a_0}{2} + \frac{1}{\sqrt{2}} \mu_A \cos \alpha$
- ▶ Bob:  $P_{\hat{\lambda}}^B(\text{yes}) = \frac{b_0}{2} - \frac{1}{\sqrt{2}} \mu_B \operatorname{sgn}(\cos \beta)$
- ▶  $\alpha$  ( $\beta$ ) is angle with pre-shared variable  $\hat{\lambda}$
- ▶  $\operatorname{sgn}(x) = +1$  ( $-1$ ) for  $x \geq 0$  ( $x < 0$ )

## Joint outcome probabilities from $M_{fb}$ and $M_{fs}$



$$P_{lhv}^{AB}(*, *) = \int \rho(\hat{\lambda}) P_{\hat{\lambda}}^A(*) P_{\hat{\lambda}}^B(*) d\hat{\lambda}$$

reproduces the singlet statistics.

- ▶  $\rho(\hat{\lambda}) = \frac{1}{4\pi}$  ( $\hat{\lambda}$  uniformly distributed over unit sphere)

## Measure of restriction on observable

By considering that observables of Alice and Bob are picked from a uniform distribution of all possible two-outcome POVMs, a measure  $r$  for % restriction on observables of any of the two parties can be defined as:

$$r = \left[ 1 - \frac{\text{Area (MOI)}}{\text{Area (POI)}} \right] \times 100$$

- ▶ In the LHV model  $\mathbb{M}_{fb}$  ( $\mathbb{M}_{fs}$ ), there is 0% (29.3%) restriction on Alice's observables where Bob's observables are restricted by 50% (29.3%)

# A general class of LHV models $\{\mathbb{M}_\kappa : \kappa \geq 0\}$

## Restriction on observable

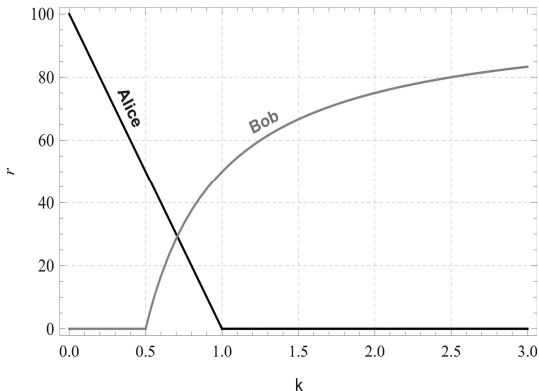
- ▶ Alice:  $\mu_A \leq \kappa \min\{a_0, 2 - a_0\}$
- ▶ Bob:  $\mu_B \leq \frac{1}{2\kappa} \min\{b_0, 2 - b_0\}$

## Model

- ▶ Alice:  $P_{\hat{\lambda}}^A(\text{yes}) = \frac{a_0}{2} + \frac{1}{2\kappa} \mu_A \cos \alpha$
- ▶ Bob:  $P_{\hat{\lambda}}^B(\text{yes}) = \frac{b_0}{2} - \kappa \mu_B \operatorname{sgn}(\cos \beta)$



## % restriction on observables for models $\mathbb{M}_\kappa$



- ▶ The subclass  $\{\mathbb{M}_\kappa : \kappa \in [1/2, 1]\}$  contains tight LHV models in as they can capture any varying degree of restrictions on Alice's and Bob's observables.

# Conclusion

- ▶ Simulation of quantum statistics for Werner state by LHV has been an interesting area for understanding the physics of entanglement
- ▶ We have studied the cases where LHV simulation is possible for singlet state.
- ▶ We find the optimal set of two outcomes observable for which singlet simulation by LHV is possible under the suggested protocol.
- ▶ It will be interesting to study whether the set can be enlarged with respect to different LHV model

# Thanks! Questions...

## References

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